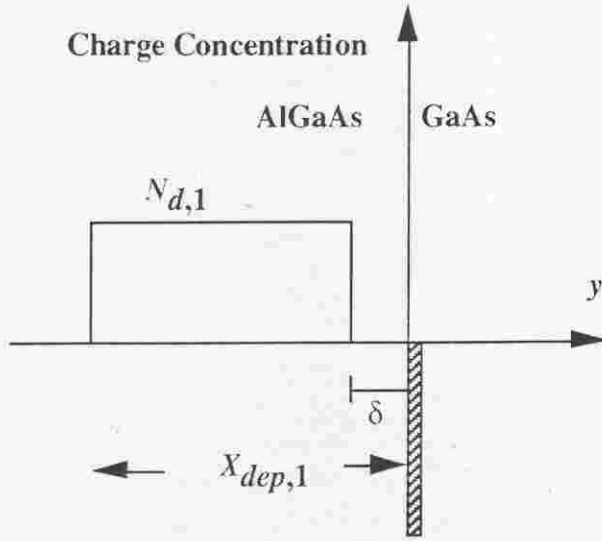
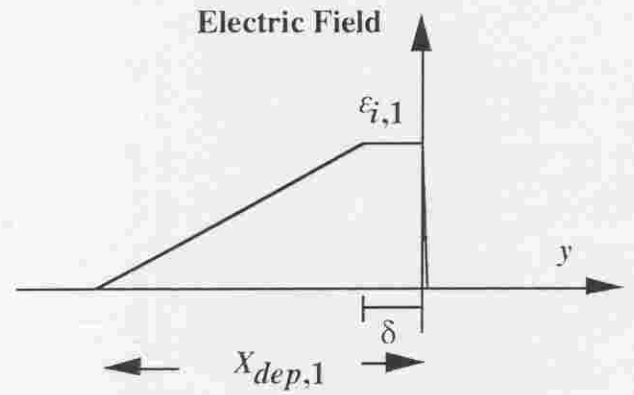


2



Charge Profile



Electric Field Profile

(b) Because we have equated  $\epsilon_{s,1} = \epsilon_{s,1}$  to  $\epsilon_s$ , the electric fields at the two sides of the interface are equal. The slopes of the electric field are 0 and  $qN_{d,1}/\epsilon_s$  in the undoped and doped portions of region 1, respectively. From an inspection of Fig. (old)5-24b, we can say,

$$\epsilon_{i,1} = \frac{q}{\epsilon_s} (X_{dep,1} - \delta) N_{d,1} \quad (1)$$

The total amount of area enclosed by the electric field is the potential drop across the two ends. Hence, we can write,

$$\epsilon_{i,1} \cdot \delta + \frac{(X_{dep,1} - \delta)}{2} \epsilon_{i,1} = \phi_{bi,1}, \quad (2)$$

where  $\phi_{bi,1}$  and  $X_{dep,1}$  are the built-in voltage and the depletion thickness in region 1, respectively. Substituting Eq. (1) into Eq. (2) eliminates the variable  $X_{dep,1}$ . The resultant equation is quadratic in  $\epsilon_{i,1}$ , which can be solved as,

$$\epsilon_s \epsilon_{i,1} = \sqrt{(qN_{d,1}\delta)^2 + 2qN_{d,1}\epsilon_s \phi_{bi,1} - qN_{d,1}\delta} \quad (3)$$

(c) Since  $\epsilon_{i,1} = \epsilon_{i,2}$ ,  $n_s$  and  $\epsilon_{i,1}$  are related through Eq. (5-105). We then arrive at the final set of equations which allow the simultaneous determination of the two unknowns:  $n_s$  and  $E_f$

$$n_s = \left[ \sqrt{(N_{d,1}\delta)^2 + \frac{2N_{d,1}\epsilon_s}{q} \phi_{bi,1}(E_f) - N_{d,1}\delta} \right] \quad (4)$$

where  $n_s$  is in  $\text{cm}^{-2}$ ; and  $E_f$  is in eV. At the room temperature,  $kT = 0.0258$  eV. For clarity, we write  $\phi_{bi,1}$  explicitly as a function of  $E_f$ . The dependence is clear from Fig. 5-21:

$$\phi_{bi,1} = \frac{\Delta E_c}{q} - \Phi_{N,1} - \frac{E_f}{q} \quad (5)$$

where  $\Phi_{N,1}$  is  $E_c - E_f$  at  $x = -\infty$ . The determination of  $\Phi_{N,1}$  can be found in Examples 1-3 and 1-4.

(d) The other equation, based on Eq. (5-98), is,

$$\ln \left[ 1 + \exp \left( \frac{E_f - 1.11 \times 10^{-9} (n_s)^{2/3}}{kT} \right) \right] + \ln \left[ 1 + \exp \left( \frac{E_f - 1.95 \times 10^{-9} (n_s)^{2/3}}{kT} \right) \right]$$

$$= \frac{n_s}{D kT} \quad (6)$$

As mentioned, the zero energy level is at the bottom tip of the triangular well. The procedure to numerically solve for  $n_s$  and  $E_f$  is to first guess a negative  $E_f$ , such that the guessed  $E_f$  is below the tip of the triangular well. From this initial guess,  $\phi_{bi,1}$  is evaluated in accordance with Eq. (5), and  $n_s$  is determined from Eq. (4). These two values of  $E_f$  and  $n_s$  are substituted into Eq. (6) to check for equality. If the two sides of Eq. (6) are different, another value of  $E_f$  is guessed. The iteration continues until a certain convergence criterion is met.

9. (a)  $E_n = 0.056 n^2$  eV. So,  $E_1 = 0.056$  eV and  $E_2 = 0.224$  eV.

$$(b) n_s = D kT \ln \left[ 1 + \exp \left( \frac{E_f - E_1}{kT} \right) \right]$$

$$(c) \frac{dn_s}{dE_f} = D kT \frac{1}{1 + \exp \left( \frac{E_f - E_1}{kT} \right)} \frac{1}{kT} \exp \left( \frac{E_f - E_1}{kT} \right)$$

$$\frac{dn_s}{dE_f} = \frac{D}{1 + \exp \left( \frac{E_1 - E_f}{kT} \right)}$$

$$\text{so, } \Delta t_b = \frac{\epsilon_s}{q^2} \frac{1 + \exp \left( \frac{E_1 - E_f}{kT} \right)}{D} = \frac{\epsilon_s}{q^2 D} \frac{\exp \left( \frac{n_s}{DkT} \right)}{\exp \left( \frac{n_s}{DkT} \right) - 1}$$

(d) when  $n_s$  is large,

$$\Delta t_b \approx \frac{\epsilon_s}{q^2 D} = \frac{1.159 \times 10^{-12}}{(1.6 \times 10^{-19})^2 \cdot 1.743 \times 10^{32}} = 2.6 \times 10^{-7} \text{ cm}$$

10. (a) It is obtained by taking the derivative of Eq. (5-151) with respect to  $x$ , evaluating the derivative at  $x = L$ , and seek the condition such that the derivative is equal to  $-\infty$  (rather than  $\epsilon_{sat}$  as in § 5-7).

$$\frac{(\alpha - 1)^2}{2\alpha} = \frac{\epsilon_{sat} L}{U_{CH}(0)}$$

(b)  $\alpha = 2 - \sqrt{3}$ , or  $U_{CH}(L) = 0.268$  V. When  $U_{CH}(L) = 0.3$  V, it is larger than 0.268 V. Hence, we expect Eq. (5-151) to hold. When  $U_{CH}(L) = 0.25$  V is smaller than 0.268 V, the saturation region has been formed and Eq. (5-151) fails.

(c) 0.642 V.

11. (a)

$$\int_0^L I_D \left( 1 - \frac{1}{\epsilon_{sat}} \frac{d}{dx} U_{CH} \right) dx = - W \mu_0 C_{ox} \int_{U_{CH}(0)}^{U_{CH}(L)} U_{CH}(x) dU_{CH}$$

$$I_D = \frac{W \mu_0 C_{ox}}{2} \times \frac{[U_{CH}^2(0) - U_{CH}^2(L)]}{L + (U_{CH}(0) - U_{CH}(L)) / \epsilon_{sat}}$$

With  $U_{CH}(0) = V_{GS} - V_T$  and  $U_{CH}(L) = V_{GS} - V_T - V_{DS}$ , we get

$$I_D = \frac{W \mu_0 C_{ox}}{L + V_{DS} / \epsilon_{sat}} \times \left[ (V_{GS} - V_T) V_{DS} - \frac{V_{DS}^2}{2} \right]$$

$$\Delta I_D^2 + B \Delta I_D - \frac{2I_{D,Sat} \epsilon_s Wh \mu_0 \epsilon_{sat} (V_{DS} - V_{D,Sat})}{\left(L + \frac{V_{D,Sat}}{\epsilon_{sat}}\right)^2} = 0$$

where  $B$  is,

$$B = \frac{2\epsilon_s Wh \mu_0 \epsilon_{sat}^2}{L + \frac{V_{D,Sat}}{\epsilon_{sat}}} - \frac{2\epsilon_s Wh \mu_0 \epsilon_{sat} (V_{DS} - V_{D,Sat})}{\left(L + \frac{V_{D,Sat}}{\epsilon_{sat}}\right)^2}$$

$$= \frac{2\epsilon_s Wh \mu_0 \epsilon_{sat}}{L + \frac{V_{D,Sat}}{\epsilon_{sat}}} \left( \epsilon_{sat} - \frac{(V_{DS} - V_{D,Sat})}{L + \frac{V_{D,Sat}}{\epsilon_{sat}}} \right)$$

Hence, the drain current as a function of  $V_{DS}$  is given by,

$$I_D = I_{D,Sat} - \frac{B}{2} + \left[ \frac{B^2}{4} + \frac{2I_{D,Sat} \epsilon_s Wh \mu_0 \epsilon_{sat} (V_{DS} - V_{D,Sat})}{\left(L + \frac{V_{D,Sat}}{\epsilon_{sat}}\right)^2} \right]^{1/2}$$

12. From Eqs. (5-99), (5-106) and (5-107),

$$n_s = D \left[ 1.95 \times 10^{-9} (n_s)^{2/3} - 1.11 \times 10^{-9} (n_s)^{2/3} \right] + 2D \left[ E_f - 1.95 \times 10^{-9} (n_s)^{2/3} \right]$$

$$\text{or, } E_f = \frac{n_s}{2D} + \frac{3.06 \times 10^{-9}}{2} n_s^{2/3}$$

when  $n_s = 2 \times 10^{12} \text{ cm}^{-2}$ , with  $D = 2.79 \times 10^{13} \text{ cm}^{-2} \cdot \text{eV}^{-1}$ ,  $E_f = 0.279 \text{ eV}$ .

13. (a) The various parameter values given in the description are:  $t_b = 280 \text{ \AA}$ ;  $\delta = 30 \text{ \AA}$ ;  $N_{d,1} = 1 \times 10^{18} \text{ cm}^{-3}$ ; and  $\phi_B = 1 \text{ eV}$ . According to the description about Eq. (5-115),  $E_{f0} = 0.0518 \text{ eV}$ . From Eq. (1-90),  $\Delta E_c$  of an  $\text{Al}_{0.35}\text{Ga}_{0.65}\text{As}/\text{GaAs}$  heterojunction is  $0.244 \text{ eV}$ . According to Eq. (5-113):

$$\phi_{00} = \frac{1.6 \times 10^{-19} \cdot 1 \times 10^{18}}{2 \cdot 1.159 \times 10^{-12}} \left( 280 \times 10^{-8} - 30 \times 10^{-8} \right)^2 = 0.431 \text{ V}$$

The threshold voltage is found from Eq. (5-119) as:

$$V_T = \phi_B + \frac{E_{f0}}{q} - \phi_{00} - \frac{\Delta E_c}{q}$$

$$= 1.0 + 0.0518 - 0.431 - 0.244 = 0.377 \text{ V}$$

(b)  $V_{D,Sat} = V_{GS} - V_T = 0.5 - 0.377 = 0.123 \text{ V}$ . Since  $V_{DS} > V_{D,Sat}$ , the transistor is in saturation and the saturation index  $\alpha = 0$ . The gate capacitance per area is found from Eq. (5-121):

$$C'_{ox} = \frac{1.159 \times 10^{-12}}{280 \times 10^{-8} + 68 \times 10^{-8}} = 3.33 \times 10^{-7} \frac{\text{F}}{\text{cm}^2}$$

The current is given by Eq. (5-133):

$$I_D = \frac{WC'_{ox} \mu_n (V_{GS} - V_T)^2}{L}$$

$$= \frac{500 \times 10^{-4} \cdot 3.33 \times 10^{-7} \cdot 6500 (0.5 - 0.377)^2}{0.25 \times 10^{-4} \cdot 2} = 0.033 \text{ A}$$

14. (a) False. The statement is true only under the d.c. condition. During transient, there is also gate current.

```

z12= -ygd/dely
z21= -ydg/dely
z22= ygg/dely
z11prime= z11 + RG + RS
z12prime= z12 + RS
z21prime= z21 + RS
z22prime= z22 + RG + RS
h21= cabs(- z21prime/z22prime)
U= cabs(z21prime - z12prime)* cabs(z21prime- z12prime)
+ /(Real(z11prime)*Real(z22prime)-Real(z12prime)*Real(z21prime))
+ /4.
write (1,*) h21, U
end

```

11. The factor  $(1 - \alpha)$  in Eq. (6-170) can be found from Eq. (5-132) as  $V_{DS}/V_{DSat}$ . Therefore, according to Eq. (6-170)

$$R_{ch} \rightarrow \frac{1}{g_m} \frac{V_{DS}}{V_{DSat}} \left[ \frac{3\alpha^3 + 15\alpha^2 + 10\alpha + 2}{10(1 + \alpha)(1 + 2\alpha)^2} \right]_{\alpha=1} = \frac{1}{6 g_m} \frac{V_{DS}}{V_{DSat}}$$

12. (a)  $R_G = \frac{1}{3} R_{SHG} \frac{W}{L} = \frac{1}{3} \cdot 0.03 \cdot \frac{1000}{0.5} = 20 \Omega$

(b)  $R_G = \frac{1}{N} \times \frac{1}{3} R_{SHG} \frac{W}{L} = \frac{1}{10} \times \frac{1}{3} \cdot 0.03 \cdot \frac{100}{0.5} = 0.2 \Omega$

(c)  $R_G = \frac{1}{N} \times \frac{1}{12} R_{SHG} \frac{W}{L} = \frac{1}{5} \times \frac{1}{12} \cdot 0.03 \cdot \frac{200}{0.5} = 0.2 \Omega$

13. From Eq. (5-117),  $\Delta t_b = 68 \text{ \AA}$ . According to Eq. (5-121), the total gate oxide capacitance is:

$$C_{ox} = \frac{13.1 \times 8.85 \times 10^{-14}}{(300 + 68) \times 10^{-8}} \cdot 0.29 \times 32 \times 10^{-8} = 2.92 \times 10^{-14} \text{ F}$$

The overlap gate-to-drain capacitance (which is identical to the gate-to-source capacitance) is:

$$C_{gd,p} = 2.92 \times 10^{-14} \times 0.1 = 2.9 \times 10^{-15} \text{ F}$$

According to Eqs. (6-16) and (6-17), the intrinsic  $C_{gg}$  and  $C_{gd}$  at  $\alpha = 0$  are:

$$C_{gg} = 2.92 \times 10^{-14} \times \left[ \frac{2}{3} \right] = 1.95 \times 10^{-14} \text{ F}$$

$$C_{gd} = 4.6 \times 10^{-14} \times [0] = 0$$

The total  $C_{gg,t}$  is the intrinsic component plus two times the overlap capacitance. It is two times because the overlap exists at both the drain and the source sides:

$$C_{gg,t} = 1.95 \times 10^{-14} + 2 \times 2.9 \times 10^{-15} = 2.53 \times 10^{-14} \text{ F}$$

The total  $C_{gd,t}$  is equal to the intrinsic component plus the overlap component:

$$C_{gd,t} = 0 + 2.9 \times 10^{-15} = 2.9 \times 10^{-15} \text{ F}$$

The gate resistance is calculated from Eq. (6-193):

$$R_G = \frac{1}{3} \cdot 5 \cdot \frac{32}{0.29} \times \frac{1}{8} = 23 \Omega$$

According to Eq. (6-235), we have,

$$\begin{aligned} \frac{1}{2\pi f_T} &= \frac{2.53 \times 10^{-14}}{0.07} + \frac{2.53 \times 10^{-14}}{0.48} (1 + 3) \cdot 0.00573 + (1 + 3) \cdot 2.9 \times 10^{-15} \\ &= 3.74 \times 10^{-13} \text{ s} \end{aligned}$$

Therefore,  $f_T = 425 \text{ GHz}$ . To find the maximum oscillation frequency, we determine the

parameter  $\Psi$  from Eq. (6-239):

$$\begin{aligned} \Psi &= (3+1) \frac{(2.53 \times 10^{-14})^2 (0.00573)^2}{(0.07)^2} + (3+1) \frac{(2.53 \times 10^{-14})(2.9 \times 10^{-15})(0.00573)}{(0.07)} \\ &\quad + \frac{(2.53 \times 10^{-14})^2 (0.00573)}{(0.07)^2} \\ &= 1.0 \times 10^{-27} \text{ F}\cdot\text{s} \end{aligned}$$

From Eq. (6-238), we then have,

$$f_{max} = \sqrt{\frac{425 \times 10^9}{8\pi \cdot 23 \cdot 2.9 \times 10^{-15} \left(1 + \frac{2\pi \cdot 425 \times 10^9}{2.9 \times 10^{-15}} 1 \times 10^{-27}\right)}}$$

So,  $f_{max}$  is 363 GHz.

14. The change in geometry modifies the gate resistance. Because  $f_T$  does not depend on the gate resistance, the cutoff frequency is still 425 GHz. However,  $f_{max}$  is affected.

$$R_G = \frac{1}{3} 5 \frac{32 \times 8}{0.29} = 1471 \Omega$$

From Eq. (6-238) and the  $\Psi$  from the solution of Problem 13:

$$f_{max} = \sqrt{\frac{425 \times 10^9}{8\pi \cdot 1471 \cdot 2.9 \times 10^{-15} \left(1 + \frac{2\pi \cdot 425 \times 10^9}{2.9 \times 10^{-15}} 1 \times 10^{-27}\right)}}$$

The maximum oscillation frequency is 45 GHz.

15. MSG is equal to the magnitude of  $y_{21}/y_{12}$ . From Eq. (2-226), we have,  $y_{21} = g_m - j\omega C_{dg,t}$  and  $y_{12} = -j\omega C_{gd,t}$ . Therefore, MSG is,

$$MSG = \left| \frac{g_m - j\omega C_{dg,t}}{-j\omega C_{gd,t}} \right| = \frac{\sqrt{g_m^2 + \omega^2 C_{dg,t}^2}}{\omega^2 C_{gd,t}^2}$$

When the parasitic capacitances are neglected, then,

$$MSG = \frac{\sqrt{g_m^2 + \omega^2 C_{dg}^2}}{\omega^2 C_{gd}^2}$$

The second expression is problematic at times because  $C_{gd}$ , the intrinsic device capacitance, is zero when the device is in saturation.  $C_{gd,t}$ , in contrast, includes the parasitic capacitance between the gate and the drain and is never zero.