Tunnel junctions



Readings this lecture covers

o Ferry pp. 91-101, 114-117o Reference: Hanson Ch. 6

Quantum tunnel probability



Band diagram for tunnel junction





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$$I = e \left(\frac{\# electrons}{\text{sec ond}} \bigg|_{R-L} - \frac{\# electrons}{\text{sec ond}} \bigg|_{L-R} \right)$$

$$\frac{\# electrons}{\text{sec ond}} \bigg|_{L-R} = \sum_{leftelectronstates rightelectronstates} \left(\Pr ob_{leftelectronstateoccupied} \right) \left(\Pr ob_{rightelectronstateempty} \right) T$$

Treat particles in left as "particle in a box" Recall our way of labeling states, and each state has energy:

$$E = \frac{\hbar^2 (\pi/L)^2}{2m} (n_x^2 + n_y^2 + n_z^2)$$

$$\frac{\# electrons}{\text{sec ond}} \bigg|_{L-R} \to \sum_{n_x, n_y, n_z} \sum_{m_x, m_y, m_z} \left(\Pr ob_{\text{leftelectronstateoccupied}} \right) \left(\Pr ob_{\text{rightelectronstateempty}} \right) T$$
$$\to \sum_{n_x, n_y, n_z} \sum_{m_x, m_y, m_z} \left(P_{n_x, n_y, n_z} \right) \left(1 - P_{m_x, m_y, m_z} \right) T$$

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$$\frac{\# electrons}{\text{sec ond}}\Big|_{L-R} = \sum_{n_x, n_y, n_z} \sum_{m_x, m_y, m_z} \left(\mathbf{P}_{n_x, n_y, n_z} \right) \left(1 - \mathbf{P}_{m_x, m_y, m_z} \right) T$$

Energy and momentum are conserved in physics so:

$$T = 0$$
 unless

 $n_{x} = m_{x}$ $n_{y} = m_{y}$ $E_{left} - eV = E_{right}$

$$\begin{split} E_{left} - eV &= E_{right} \\ \Rightarrow \frac{\hbar^2 (\pi/L)^2}{2m} \left(n_x^2 + n_y^2 + n_z^2 \right) - eV = \frac{\hbar^2 (\pi/L)^2}{2m} \left(m_x^2 + m_y^2 + m_z^2 \right) \\ \Rightarrow \frac{\hbar^2 (\pi/L)^2}{2m} n_z^2 - eV = \frac{\hbar^2 (\pi/L)^2}{2m} m_z^2 \end{split}$$



 $\frac{\# electrons}{\text{sec ond}} \bigg|_{L-R} \to \sum_{n_x, n_y, n_z} (f(E_L)) (1 - f(E_L + eV)) T$ A similar calculation shows: $\frac{\# electrons}{\text{sec ond}} \bigg|_{R-L} \to \sum_{n_x, n_y, n_z} (f(E_L + eV)) (1 - f(E_L)) T$

Since:

$$I = e \left(\frac{\# electrons}{\text{sec ond}} \bigg|_{R-L} - \frac{\# electrons}{\text{sec ond}} \bigg|_{L-R} \right)$$

We have:

$$I = e \sum_{n_x, n_y, n_z} \left[\left(f(E_L) - f(E_L + eV) \right) \right] T$$

A nice, simple result.

$$I = e \sum_{n_x, n_y, n_z} \left[\left(f(E_L) - f(E_L + eV) \right) \right] T$$
$$I = e \sum_{n_x, n_y} \sum_{n_z} \left[\left(f(E_L) - f(E_L + eV) \right) \right] T$$

In the macro world, states are very finely spaced and we have (discuss): *(Later in the class we will see that this fails in nanosized circuits.)*

$$\sum_{n_x} \rightarrow \int dn_x \quad \sum_{n_y} \rightarrow \int dn_y \quad \sum_{n_z} \rightarrow \int dn_z$$
$$H \rightarrow e \int dn_x \int dn_y \int dn_z \Big[\Big(f(E_L) - f(E_L + eV) \Big) \Big] T$$

$$I \rightarrow e \int dn_x \int dn_y \int dn_z \left[\left(f(E_L) - f(E_L + eV) \right) \right] T$$

$$I \rightarrow e \int dn_x \int dn_y \int \frac{m}{\hbar^2 (\pi/L)^2} \frac{1}{\sqrt{E_L - \frac{\hbar^2 (\pi/L)^2}{2m} \left(n_x^2 + n_y^2\right)}} dE_L \left[\left(f(E_L) - f(E_L + eV) \right) \right] T$$

$$I \approx e \int dn_{x} \int dn_{y} \frac{m}{\hbar^{2} (\pi/L)^{2}} \frac{1}{\sqrt{E_{F} - \frac{\hbar^{2} (\pi/L)^{2}}{2m} \left(n_{x}^{2} + n_{y}^{2}\right)}} T \int dE_{L} \Big[\Big(f(E_{L}) - f(E_{L} + eV) \Big) \Big]$$

 $\int dE_{L} \Big[\Big(f(E_{L}) - f(E_{L} + eV) \Big) \Big] \approx eV \quad \text{(show on board)}$

$$I \approx (eV) eT \frac{m}{\hbar^2 (\pi/L)^2} \int_0^\infty dn_x \int_0^\infty dn_y \frac{1}{\sqrt{E_F - \frac{\hbar^2 (\pi/L)^2}{2m} (n_x^2 + n_y^2)}}}$$
$$\boxed{I \approx (eV) (\text{constant})}$$